

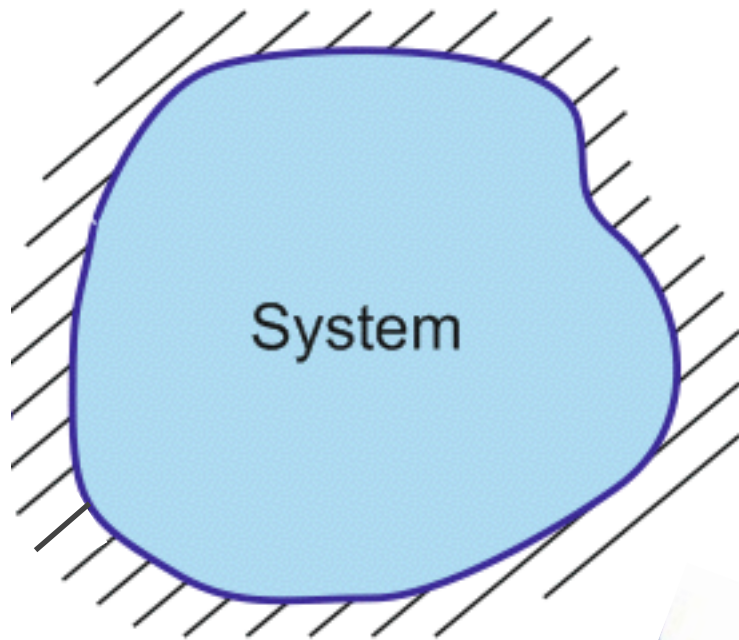
Heating in periodically driven Floquet systems

Anushya Chandran

Boston University

Floquet system

Periodically driven isolated system



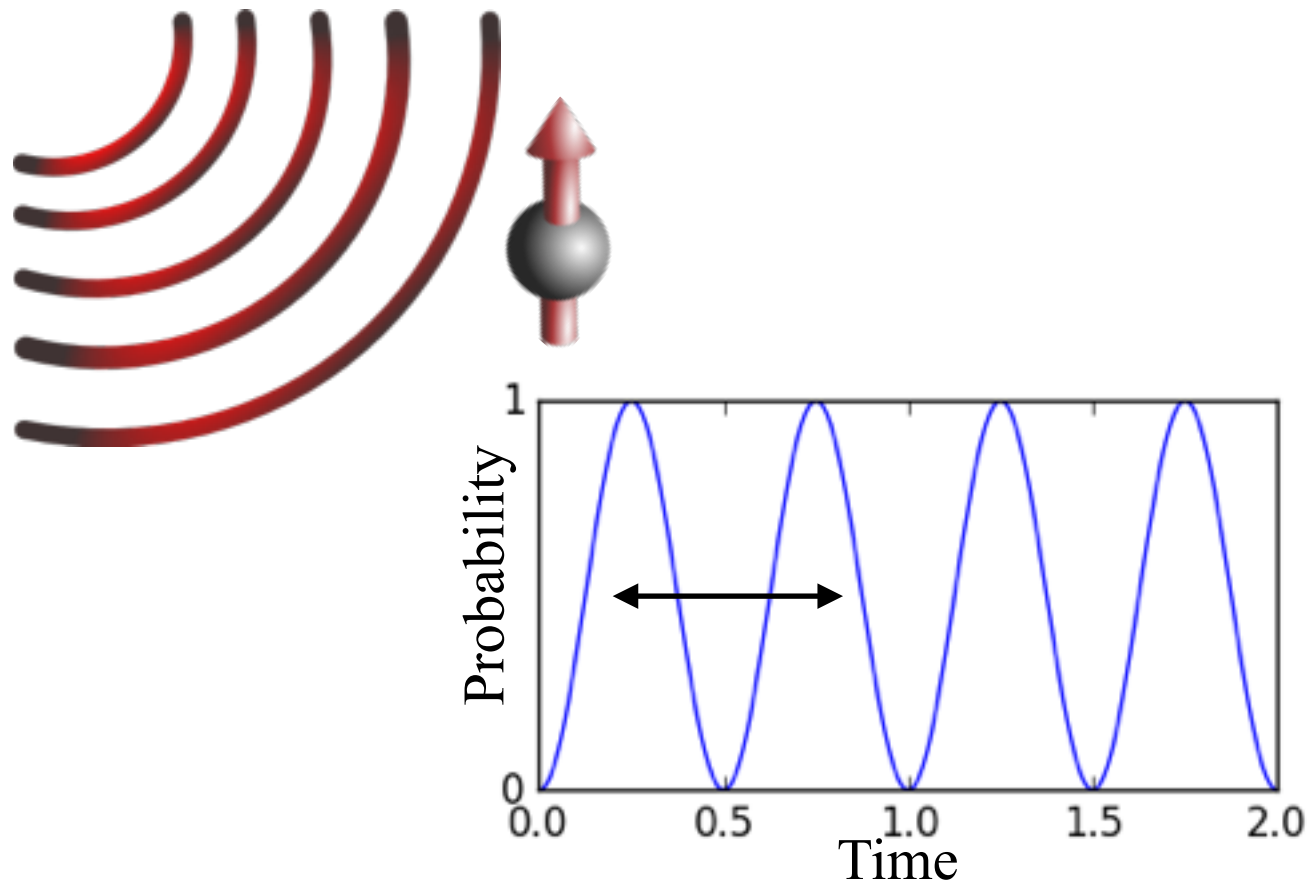
Hamiltonian H_0

$$H(t) = H_0 + V \cos(\omega t) H_1$$



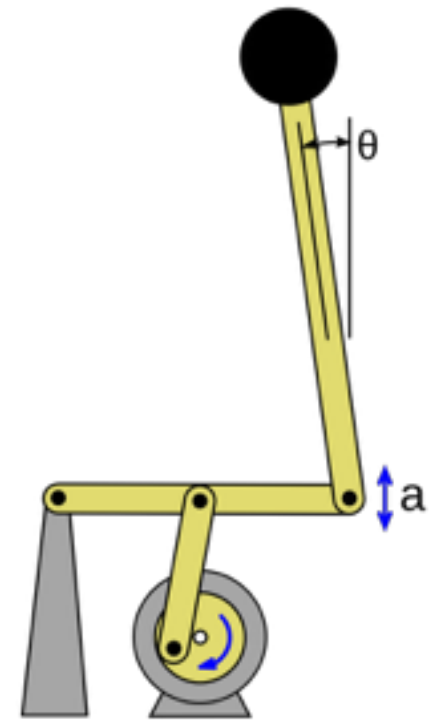
Few-body Floquet systems

Rabi oscillations



Amplitude of drive \Rightarrow Frequency

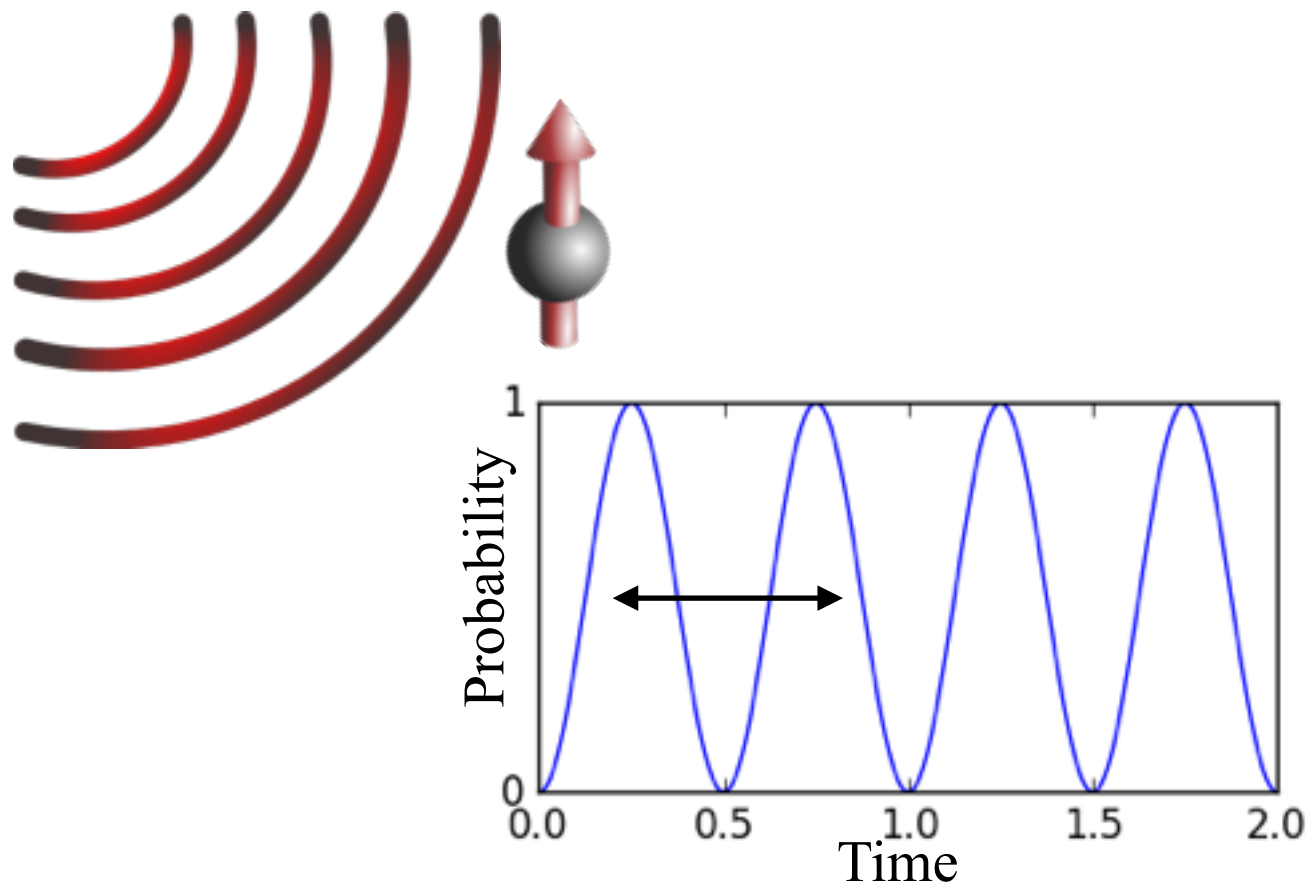
Kapitza pendulum



New stable equilibrium

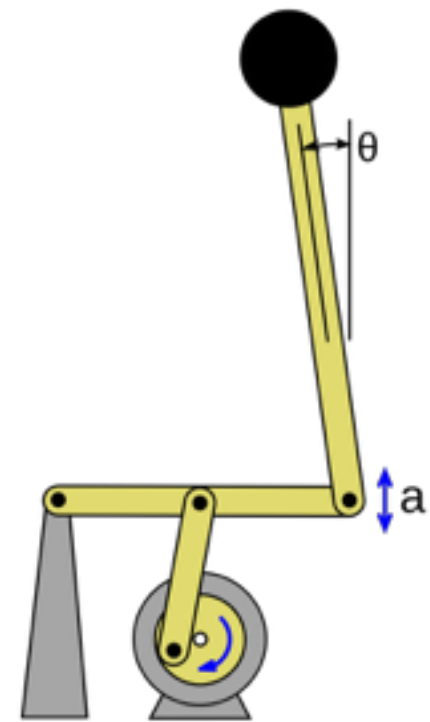
Few-body Floquet systems

Rabi oscillations



Amplitude of drive \Rightarrow Frequency

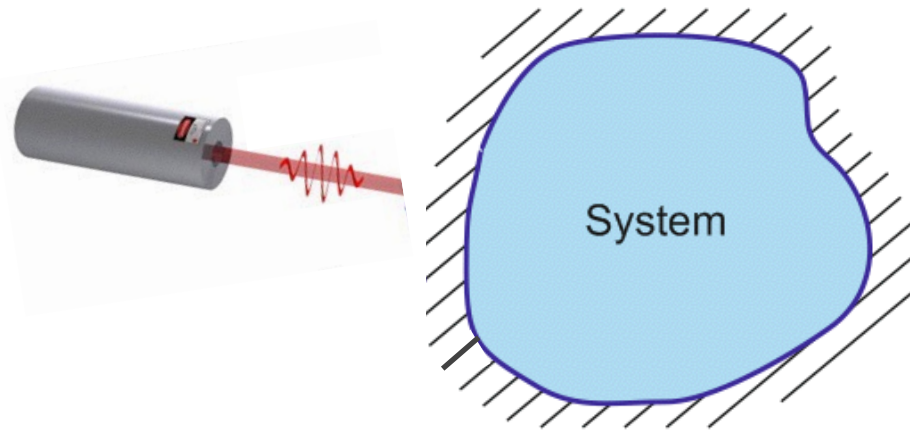
Kapitza pendulum



New stable equilibrium

Many-body?

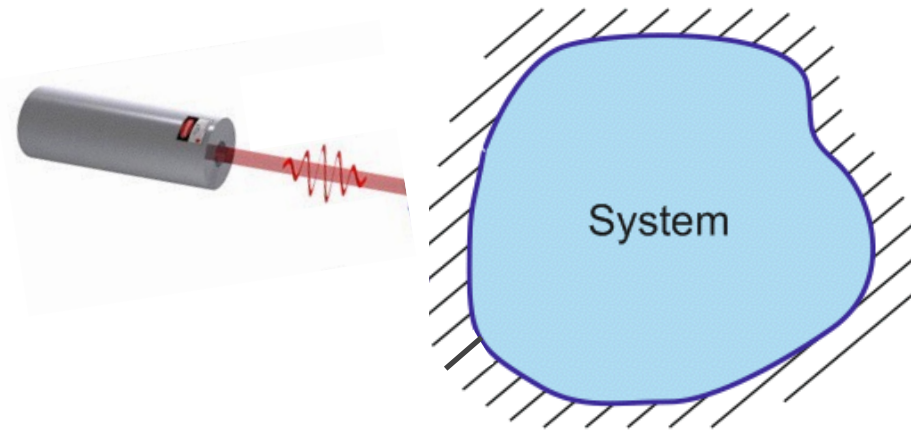
Interest: fundamental & engineering



Simplest non-equilibrium setting:
what can happen?

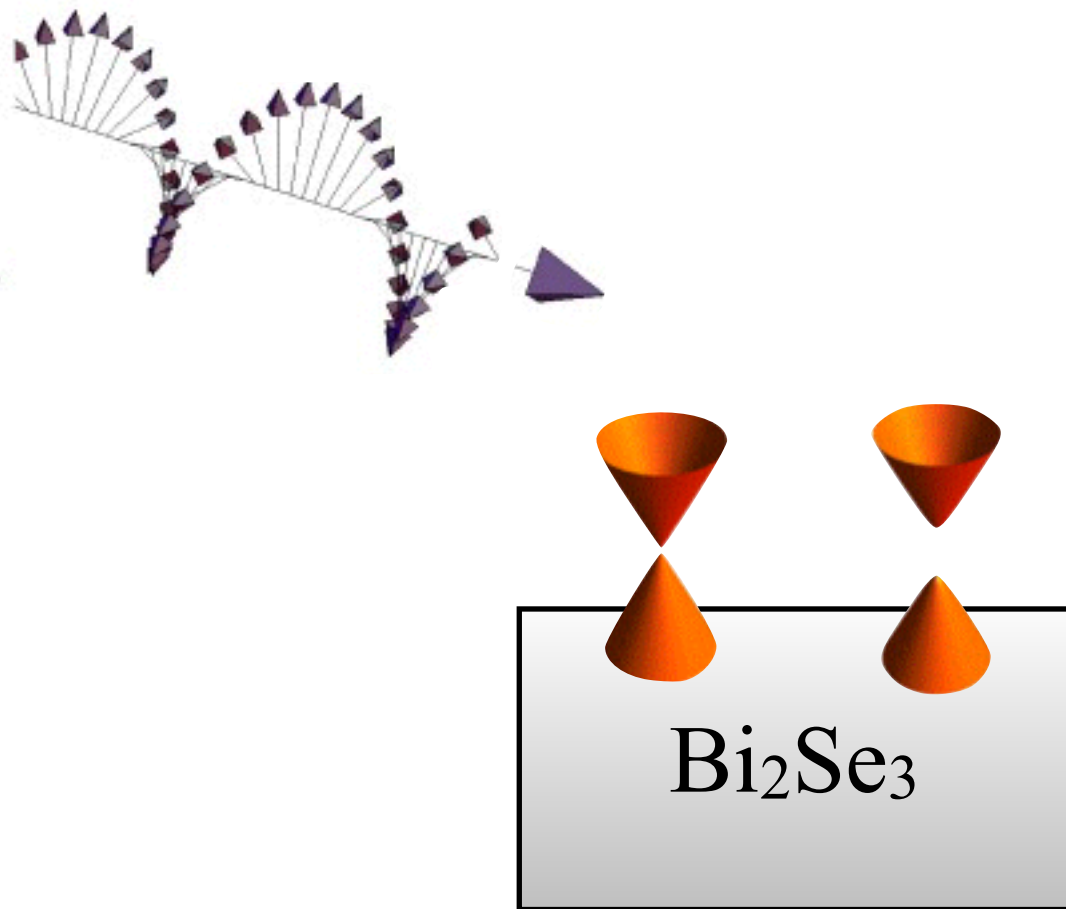
Engineer new states out of
equilibrium?

Interest: fundamental & engineering

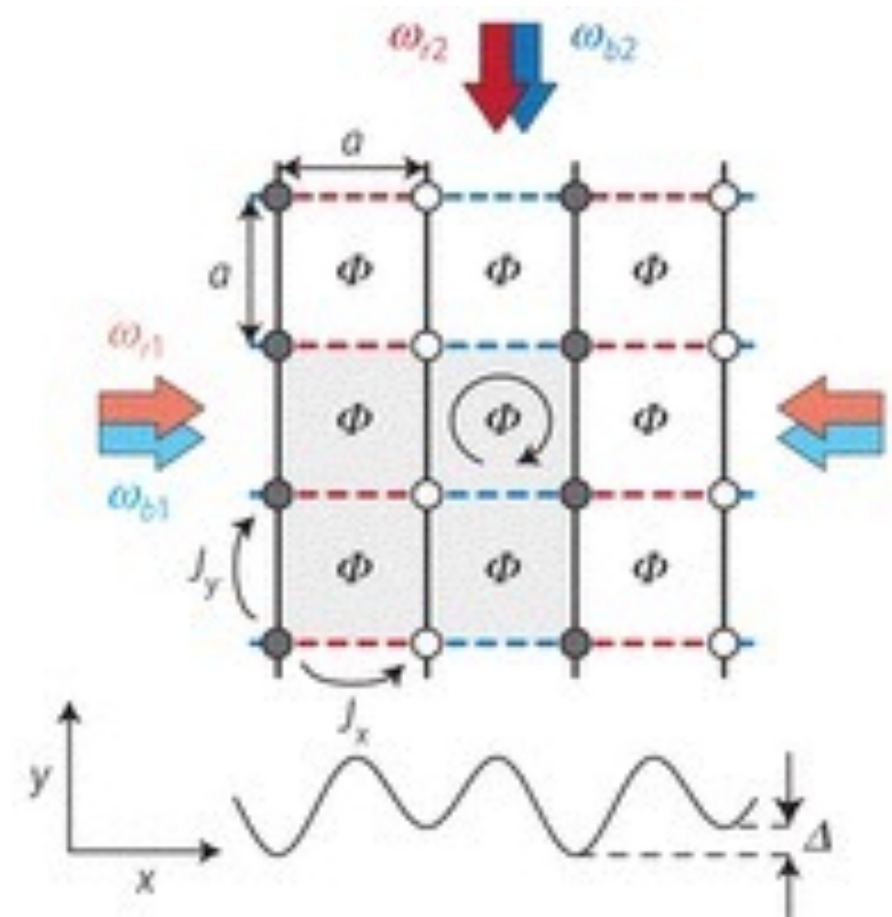


Simplest non-equilibrium setting:
what can happen?

Engineer new states out of
equilibrium?



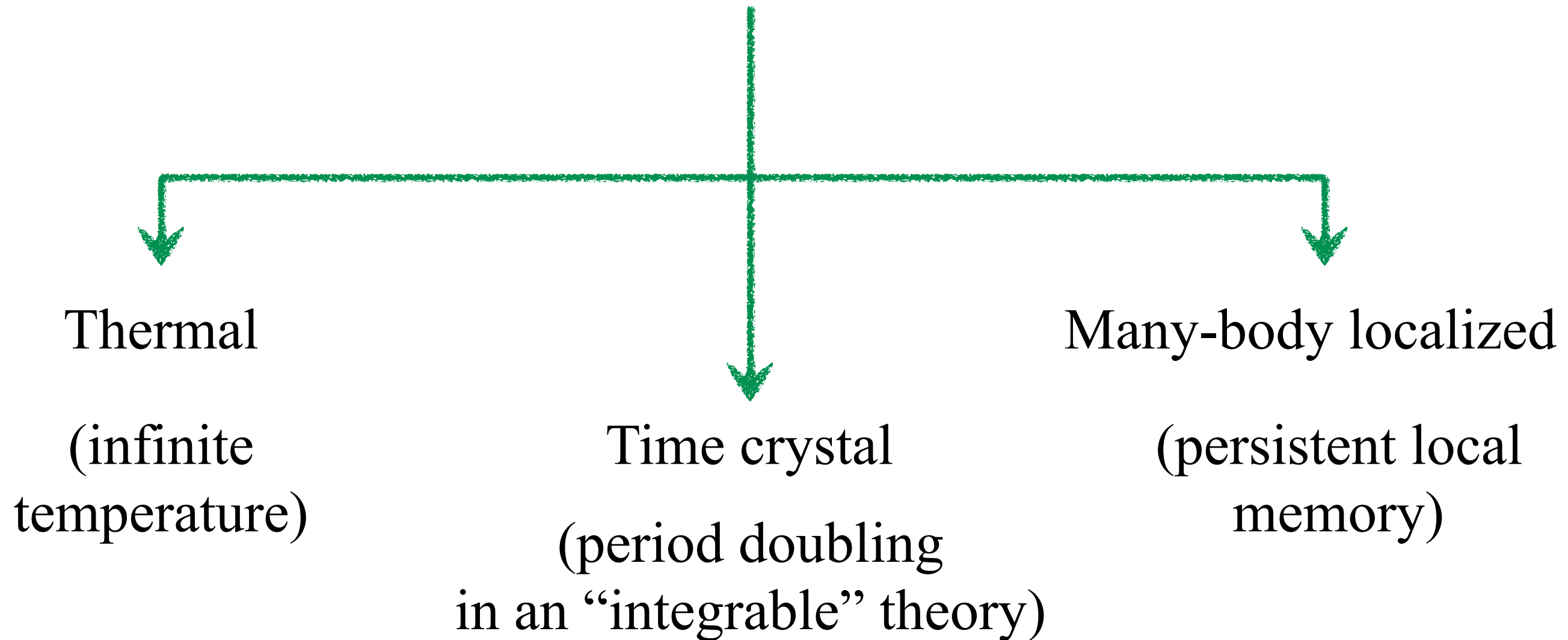
Wang et al (Gedik group)
Science (2013)



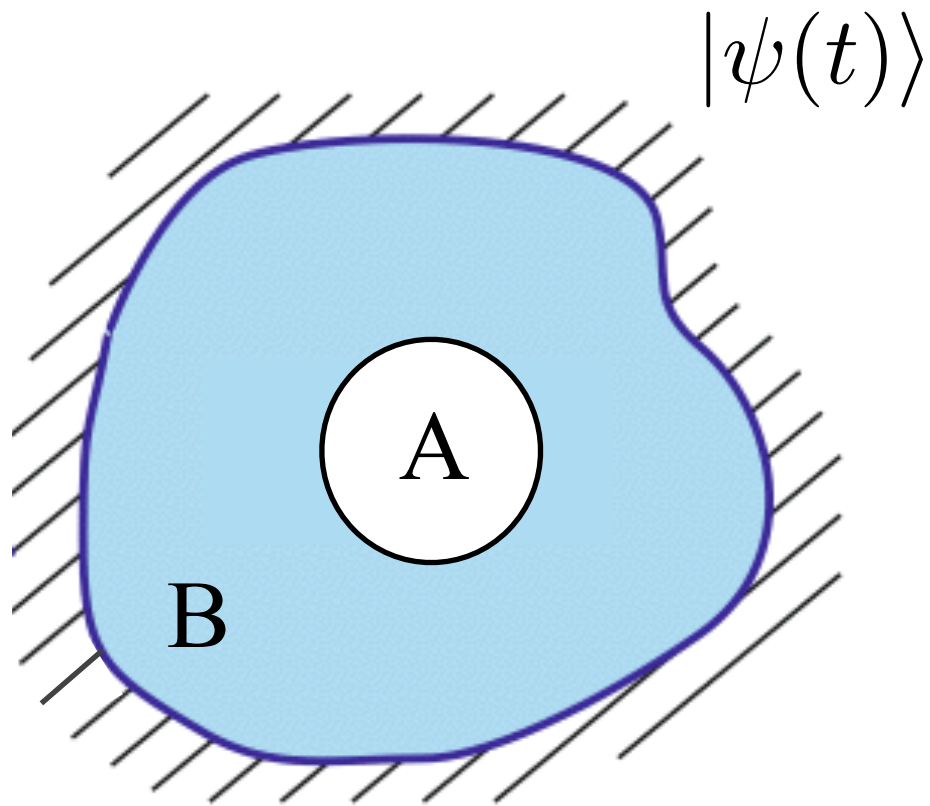
Aidelsburger et al (Bloch group)
Nature (2014)

Outline

Steady states of Floquet systems



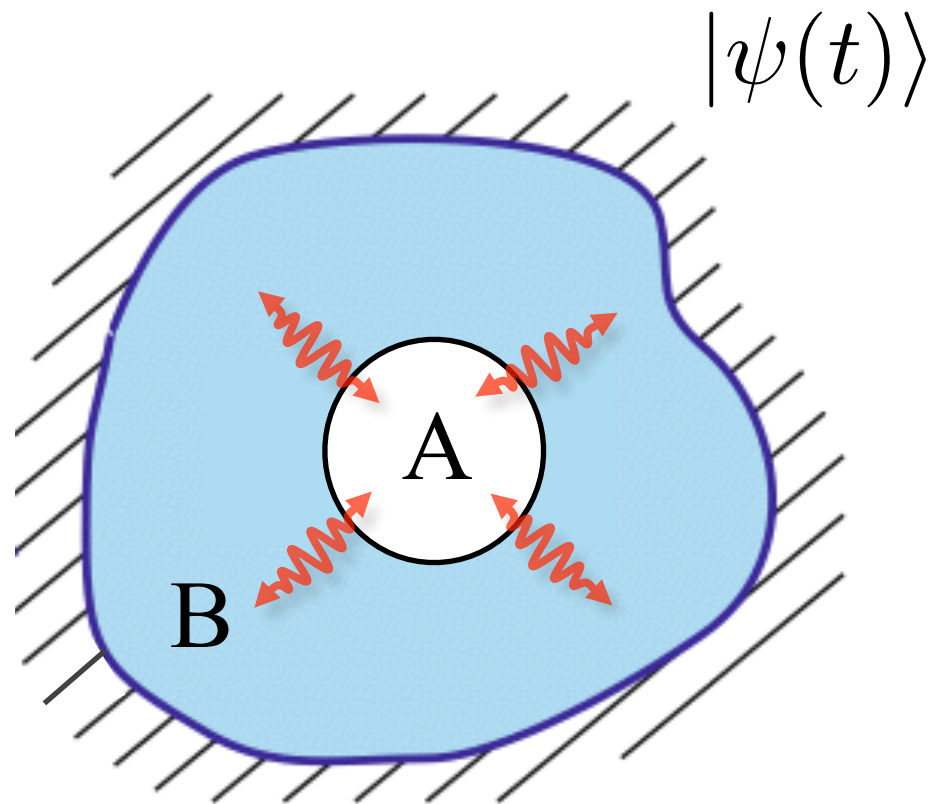
Thermalization in isolated systems



Local state

$$\rho_A(t) = \text{Tr}_B |\psi(t)\rangle \langle \psi(t)|$$

Thermalization in isolated systems



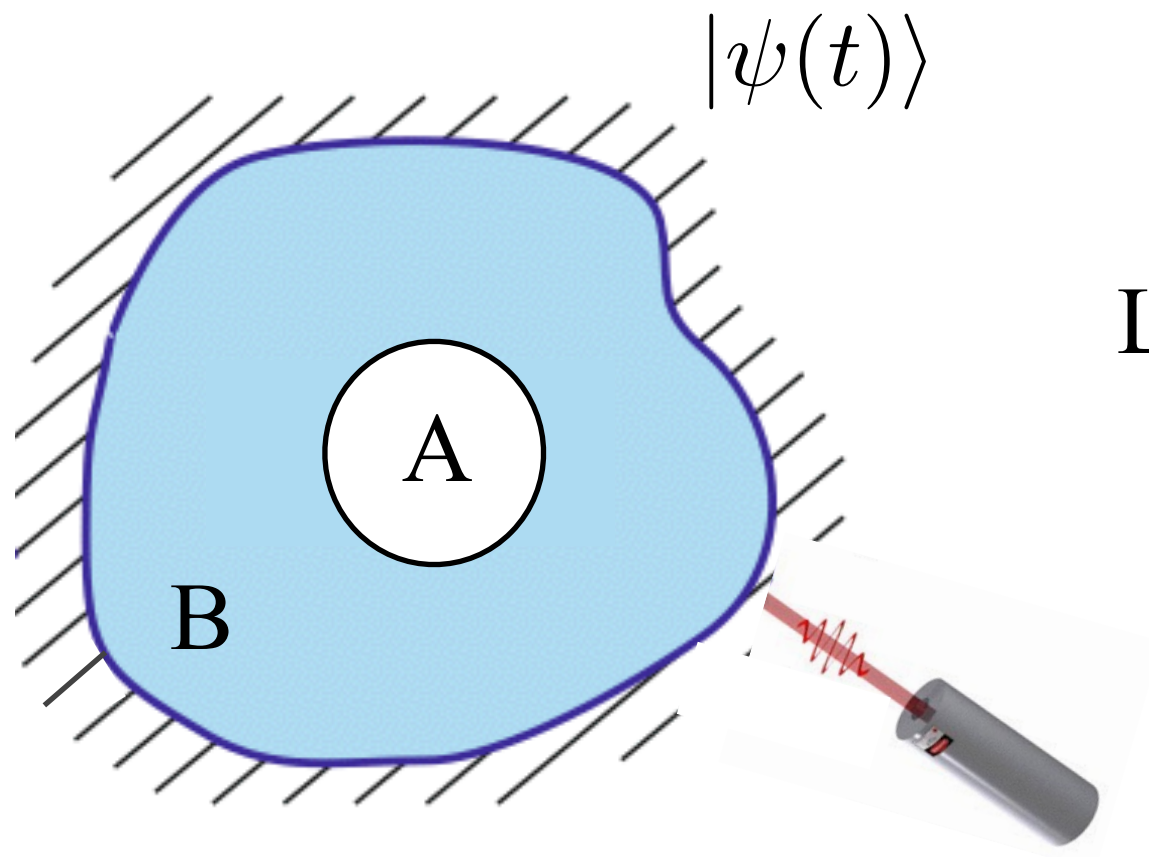
Local state

$$\rho_A(t) = \text{Tr}_B |\psi(t)\rangle \langle \psi(t)|$$

No driving

$$\lim_{t \rightarrow \infty} \rho_A(t) = \frac{1}{Z} \text{Tr}_B e^{-\beta H}$$

Thermalization in isolated systems



Local state

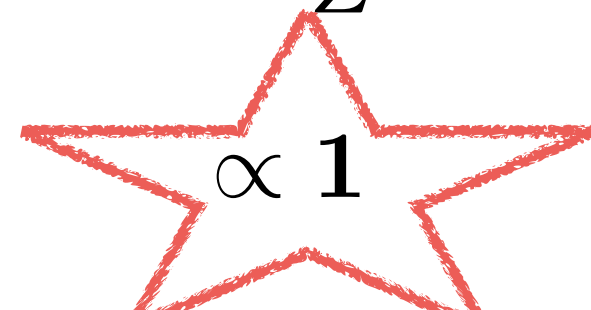
$$\rho_A(t) = \text{Tr}_B |\psi(t)\rangle \langle \psi(t)|$$

No driving

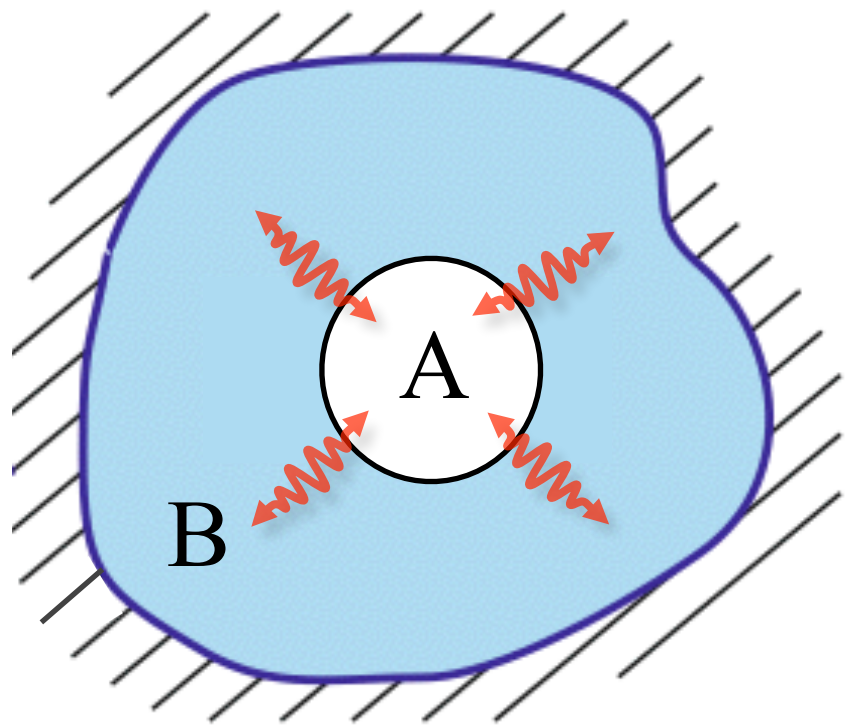
$$\lim_{t \rightarrow \infty} \rho_A(t) = \frac{1}{Z} \text{Tr}_B e^{-\beta H}$$

With driving

$$\lim_{t \rightarrow \infty} \rho_A(t) = \frac{1}{Z} \text{Tr}_B e^{-\beta H}$$



Eigenstate thermalization hypothesis (ETH)



For all eigenstates E_i
at inverse temperature β

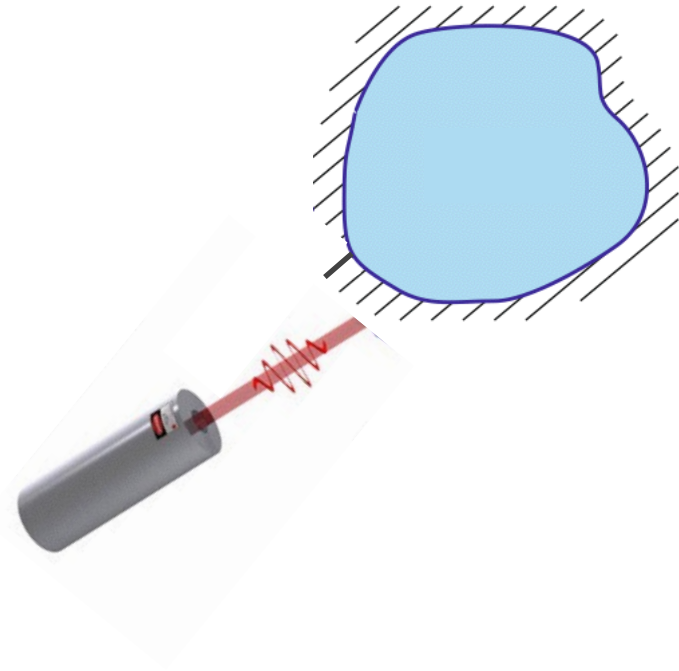
$$\rho_A = \text{Tr}_B |E_i\rangle \langle E_i| = \frac{1}{Z} \text{Tr}_B e^{-\beta H}$$

$$H|E_i\rangle = E_i|E_i\rangle$$

ETH \Rightarrow thermalization

Generically thermalization seems to \Rightarrow ETH

Driven eigenstates

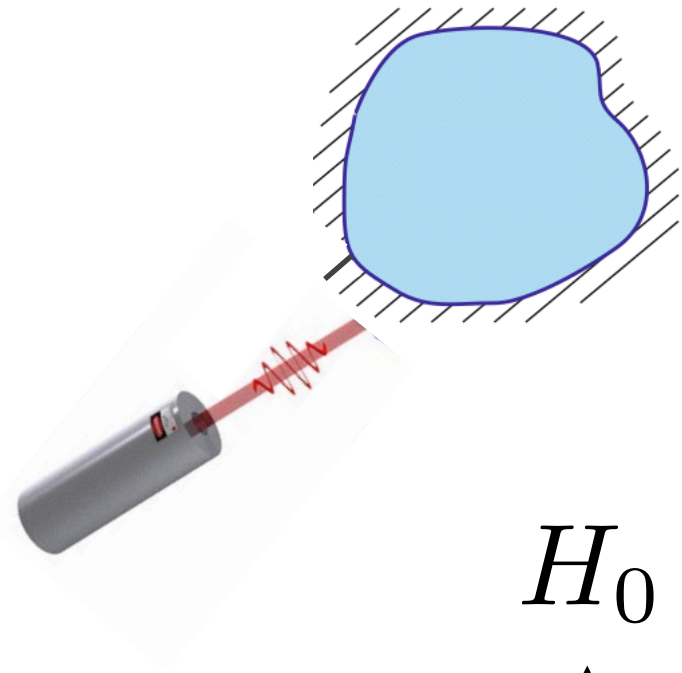


$$H(t) = H_0 + V \cos(\omega t) H_1$$

Floquet/periodic
evolution:

$$U(T) = T e^{-i \int_0^T H(t) dt'}$$

Driven eigenstates

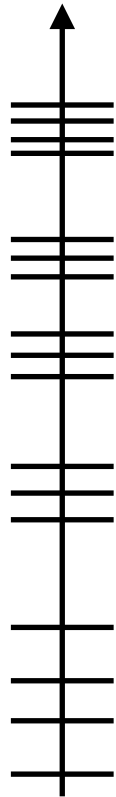


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Floquet/periodic
evolution:

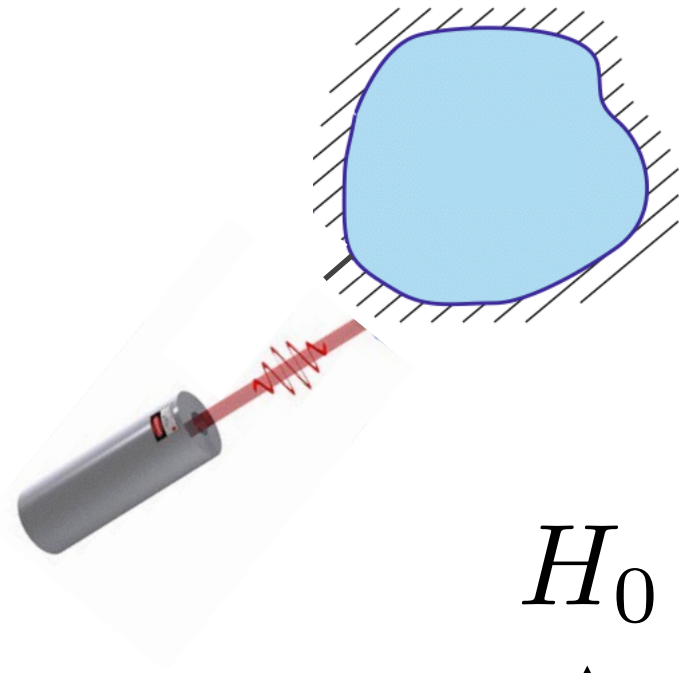
$$U(T) = T e^{-i \int_0^T H(t) dt'}$$

H_0



Energy

Driven eigenstates

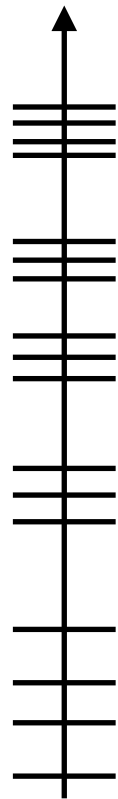


$$H(t) = H_0 + V \cos(\omega t) H_1$$

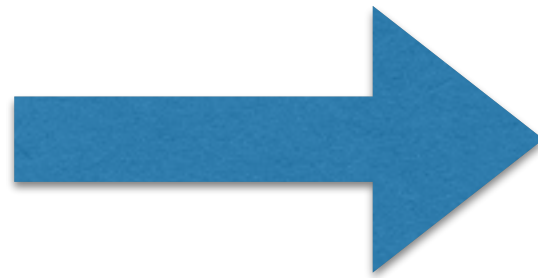
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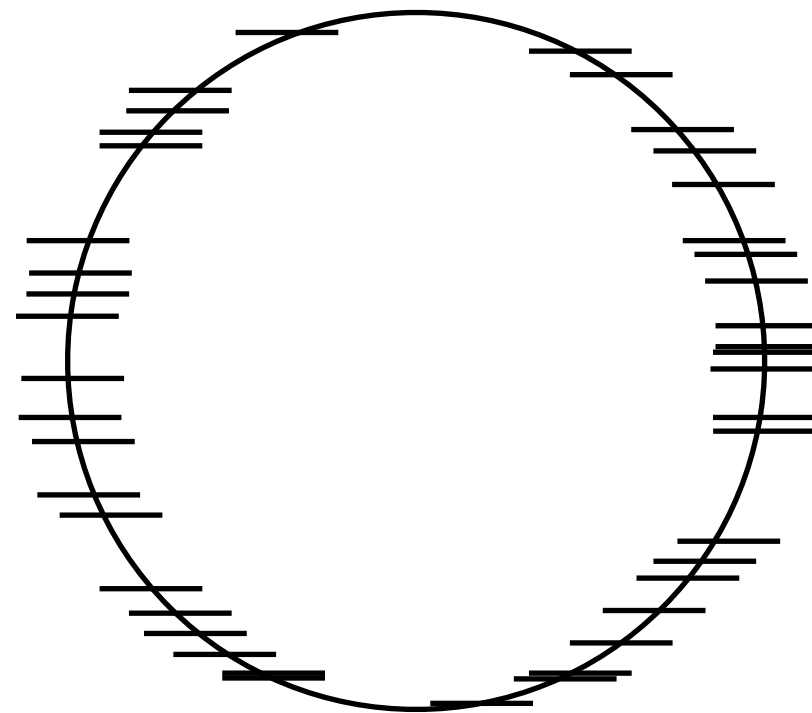
H_0



Energy



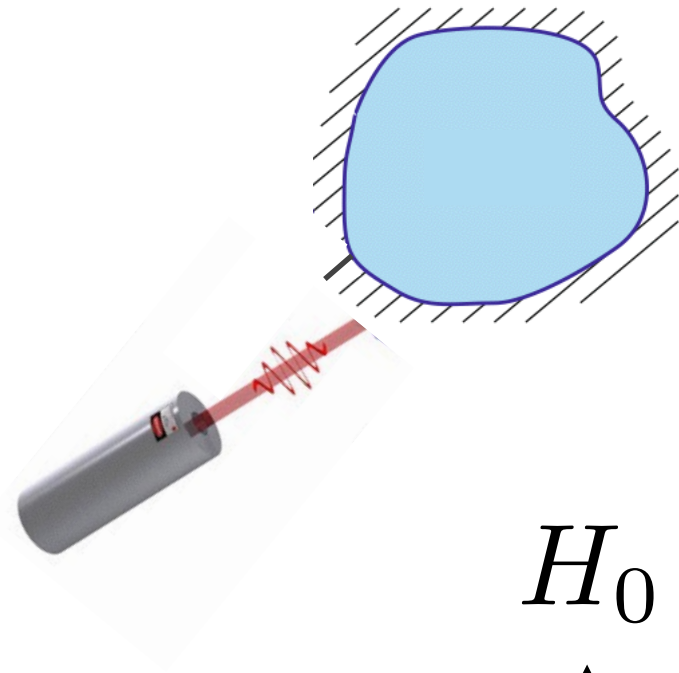
$U(T)$



Floquet Quasi-energy

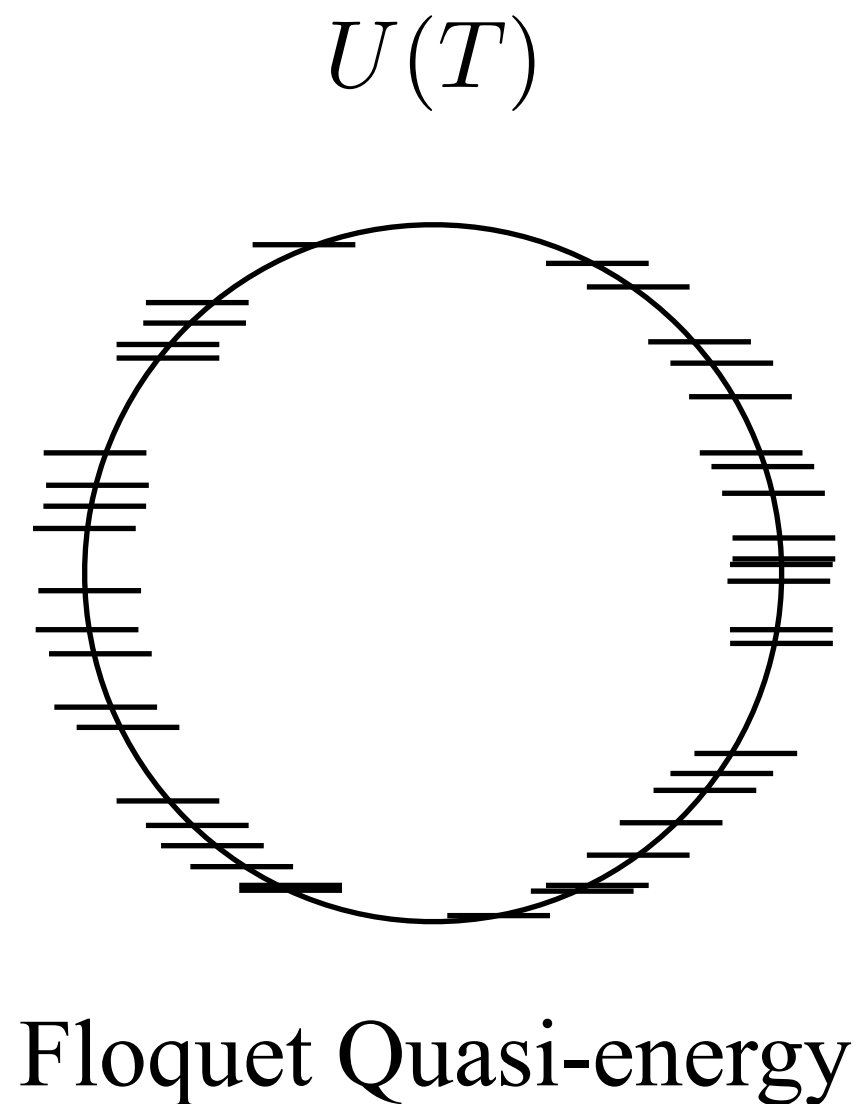
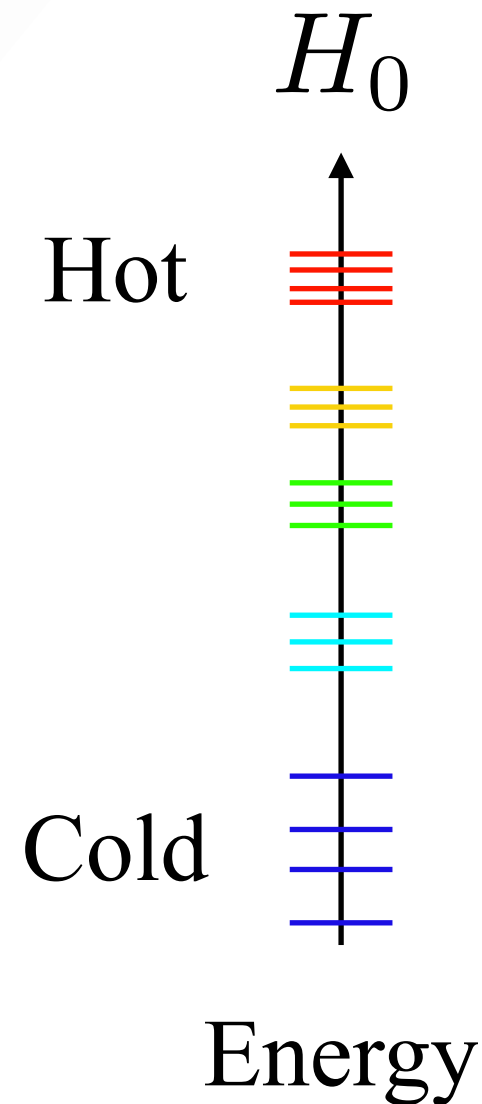
Driven eigenstates

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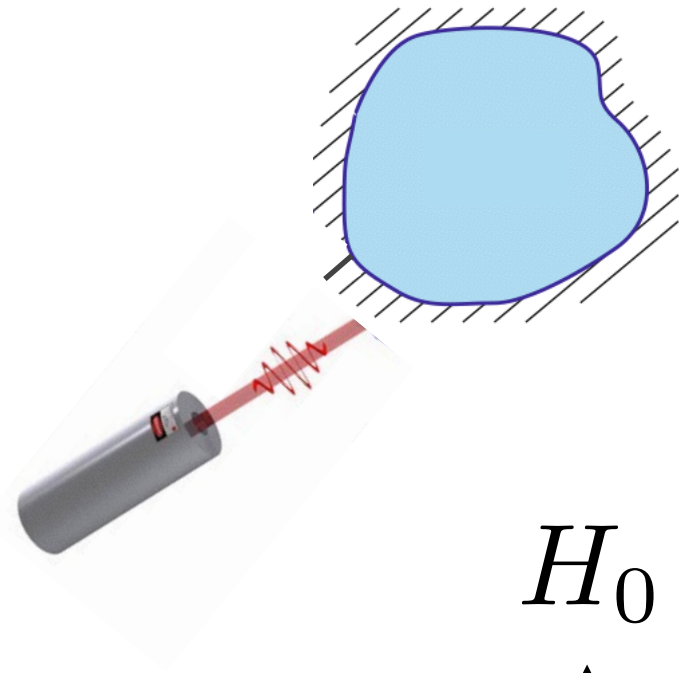


$$H(t) = H_0 + V \cos(\omega t) H_1$$

Undriven eigenstates ETH



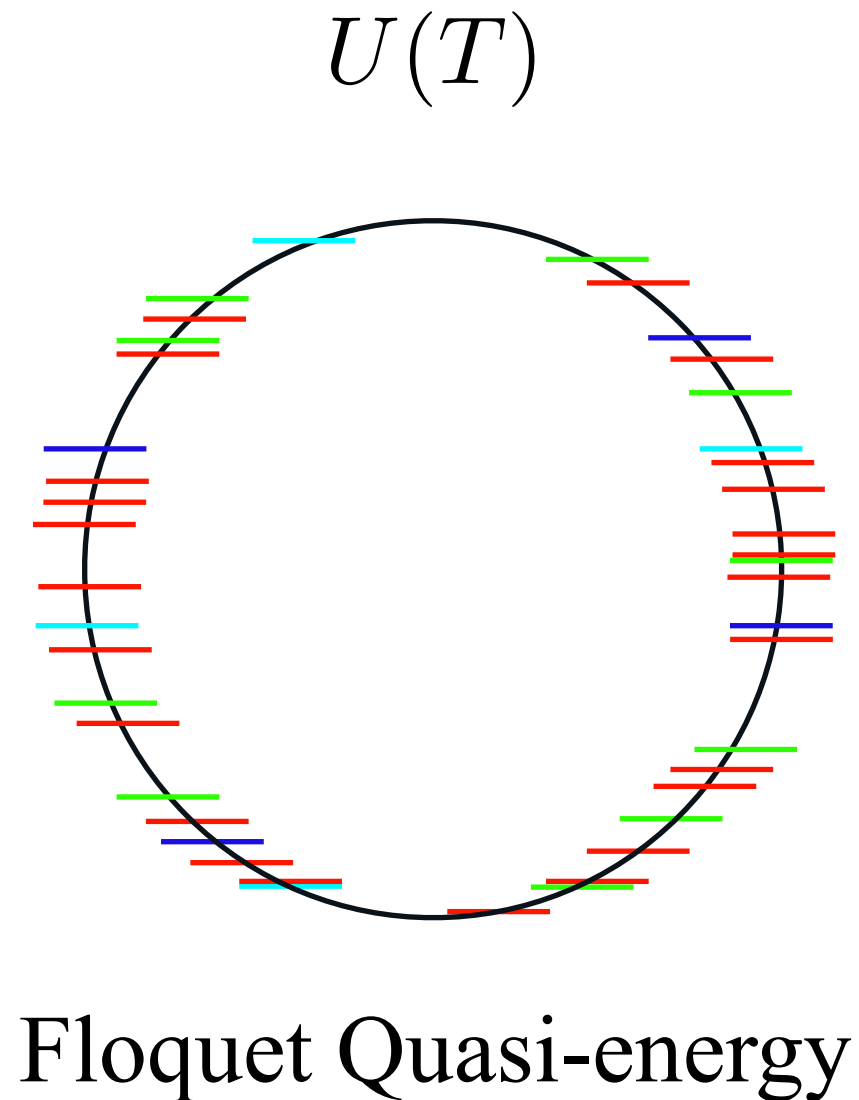
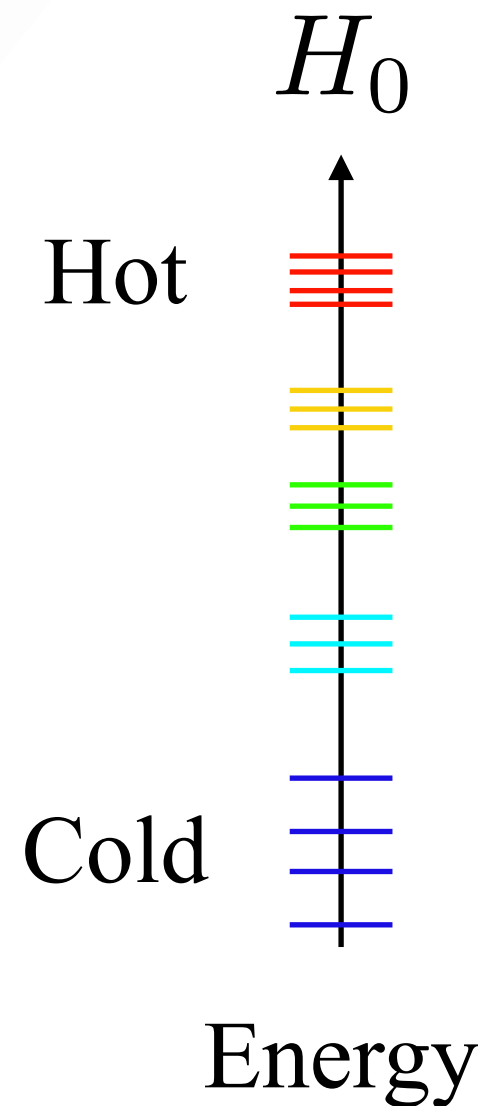
# Driven eigenstates



$$H(t) = H_0 + V \cos(\omega t) H_1$$

Undriven eigenstates ETH

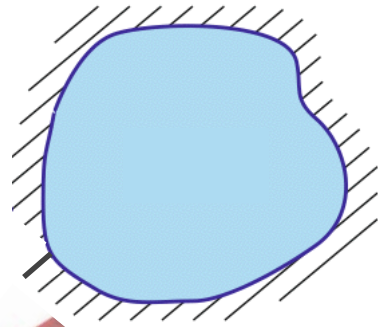
Driven eigenstates ?





# Driven eigenstates

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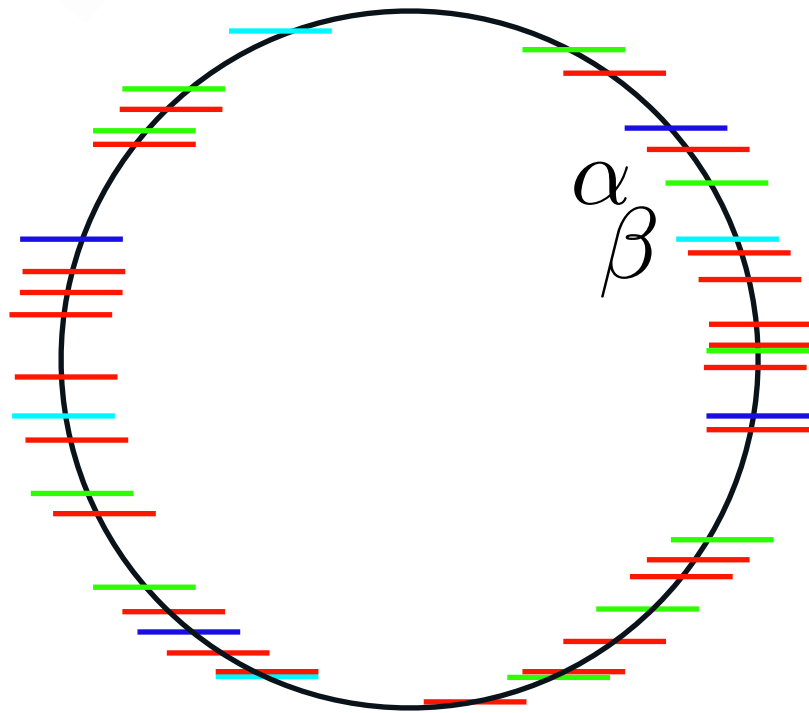
Undriven eigenstates ETH

$$H(t) = H_0 + V \cos(\omega t) H_1$$

Driven eigenstates ?

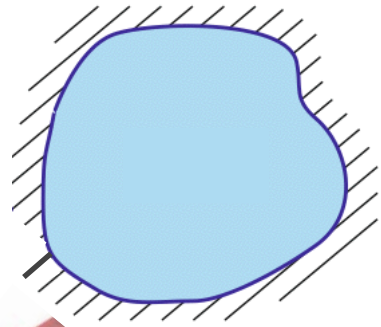
$U(T)$

Local drive: $\langle E_\beta | U(T) | E_\alpha \rangle \sim \frac{1}{\sqrt{2^L}}$



Floquet Quasi-energy

Driven eigenstates



Undriven eigenstates ETH

$$H(t) = H_0 + V \cos(\omega t) H_1$$

Driven eigenstates ?

$U(T)$

Local drive:

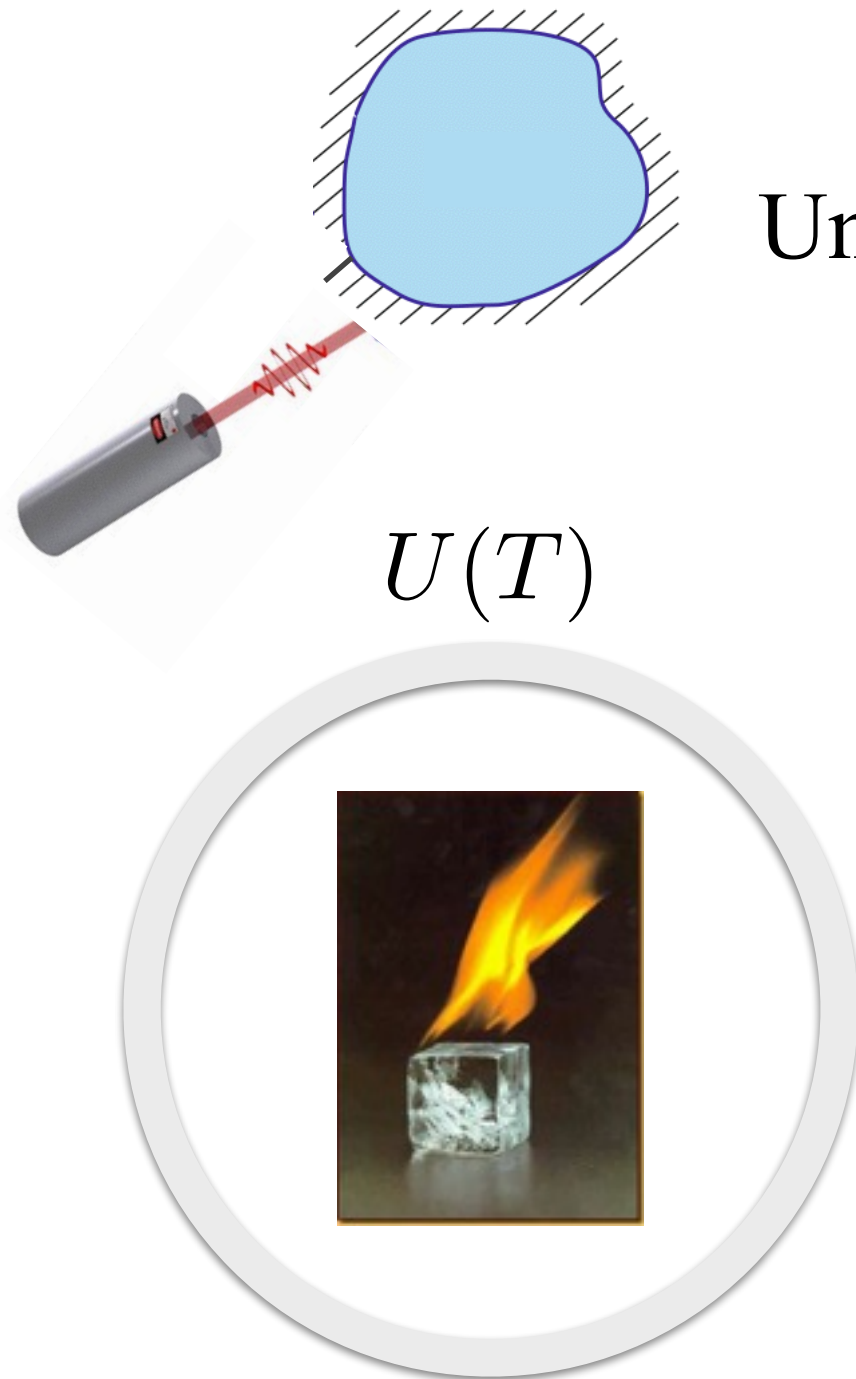
$$\langle E_\beta | U(T) | E_\alpha \rangle \sim \frac{1}{\sqrt{2^L}}$$

α
 β $\Delta_{\alpha\beta}$

$$\Delta_{\alpha\beta} \sim \frac{1}{2^L}$$

Floquet Quasi-energy

Driven eigenstates



Undriven eigenstates ETH

Driven eigenstates ?

Local drive: $\langle E_\beta | U(T) | E_\alpha \rangle \sim \frac{1}{\sqrt{2^L}}$

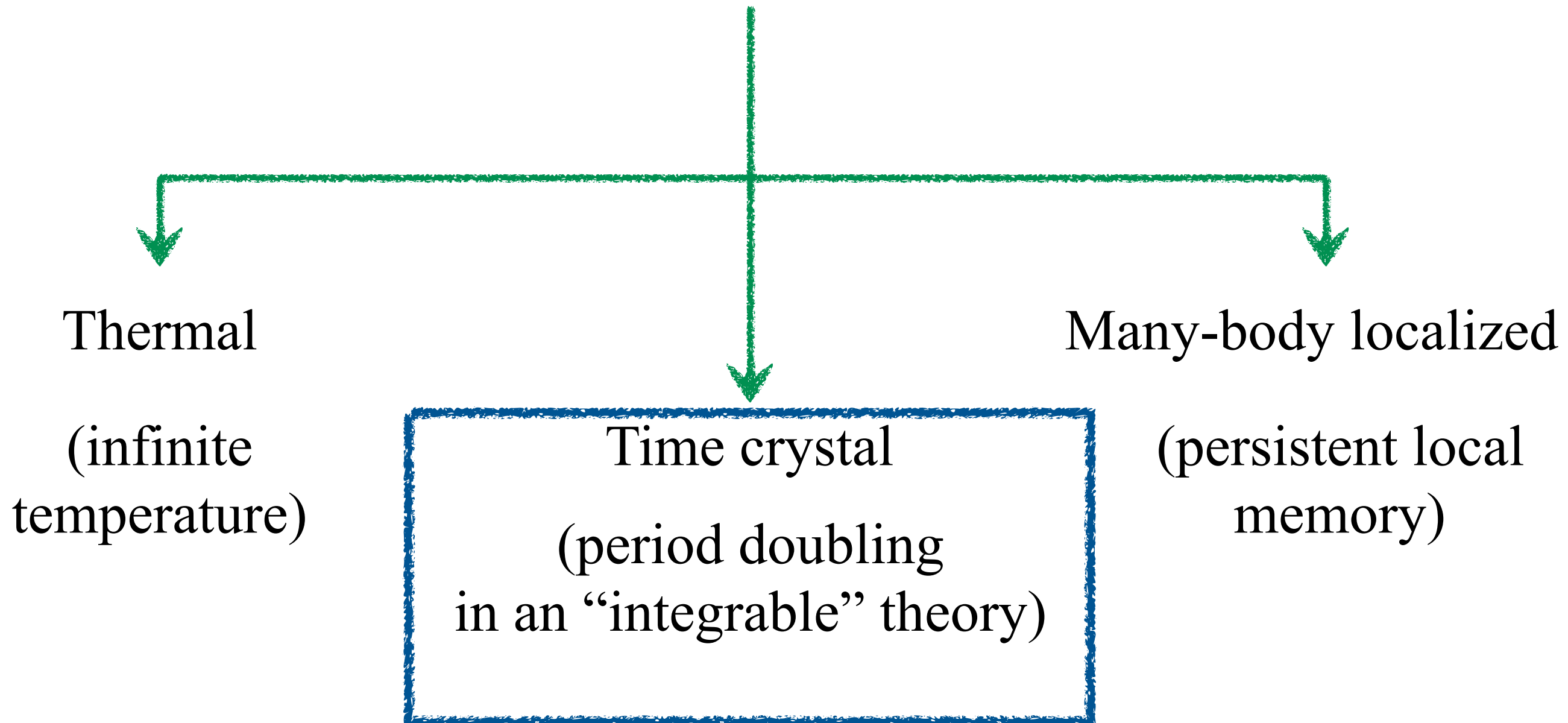
$$\Delta_{\alpha\beta} \sim \frac{1}{2^L}$$

Floquet eigenstates mix all temperatures!

Floquet Quasi-energy

Outline

Steady states of Floquet systems



Interacting driven bosons

Driven $O(N)$ model

$$H(t) = \frac{1}{2} \int d^d x (|\Pi|^2 + |\nabla \Phi|^2 + r(t) |\Phi|^2 + \lambda |\Phi|^4)$$

$$r(t) = r_0 - r_1 \cos(\gamma t)$$

Interacting driven bosons

Driven $O(N)$ model

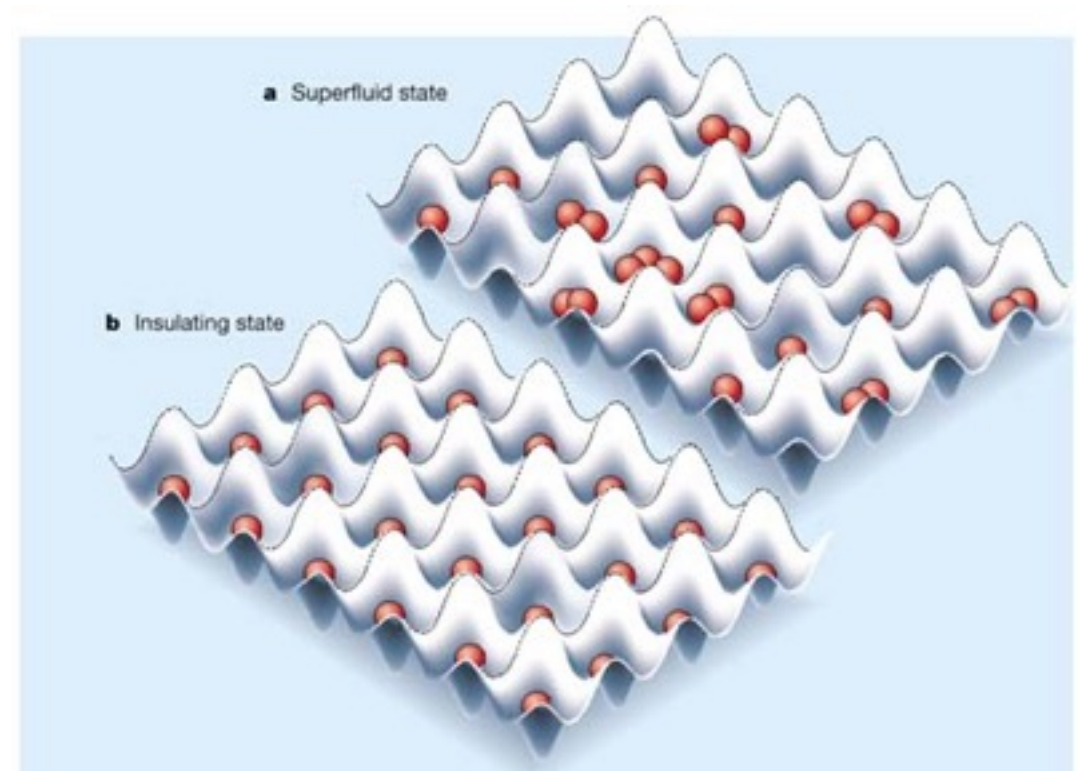
$$H(t) = \frac{1}{2} \int d^d x (|\Pi|^2 + |\nabla \Phi|^2 + r(t) |\Phi|^2 + \lambda |\Phi|^4)$$

$$r(t) = r_0 - r_1 \cos(\gamma t)$$

$O(2)$ version:

Near transition from Mott
insulator to superfluid

$r(t)$: modulating tunneling



Interacting driven bosons

Driven O(N) model

$$H(t) = \frac{1}{2} \int d^d x (|\Pi|^2 + |\nabla \Phi|^2 + r(t) |\Phi|^2 + \lambda |\Phi|^4)$$

$$r(t) = r_0 - r_1 \cos(\gamma t)$$

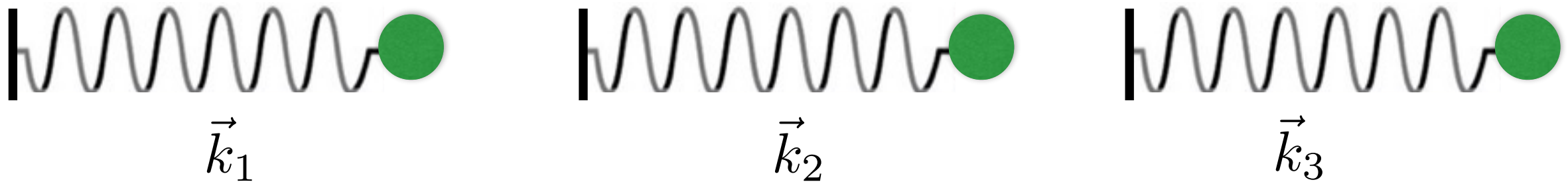
Equilibrium: canonical model for symmetry-breaking

Analytical control in the large-N limit

Self-consistent **classical** equations

$$\left(\frac{d^2}{dt^2} + |\vec{k}|^2 + r(t) + N \lambda \int^\Lambda \frac{d^d k}{(2\pi)^d} |f_{\vec{k}}(t)|^2 \right) f_{\vec{k}}(t) = 0$$

Interacting driven bosons



$$\omega_{\vec{k}}(t)^2 = |k|^2 + r(t) + r_{eff}(t)$$

Feedback term

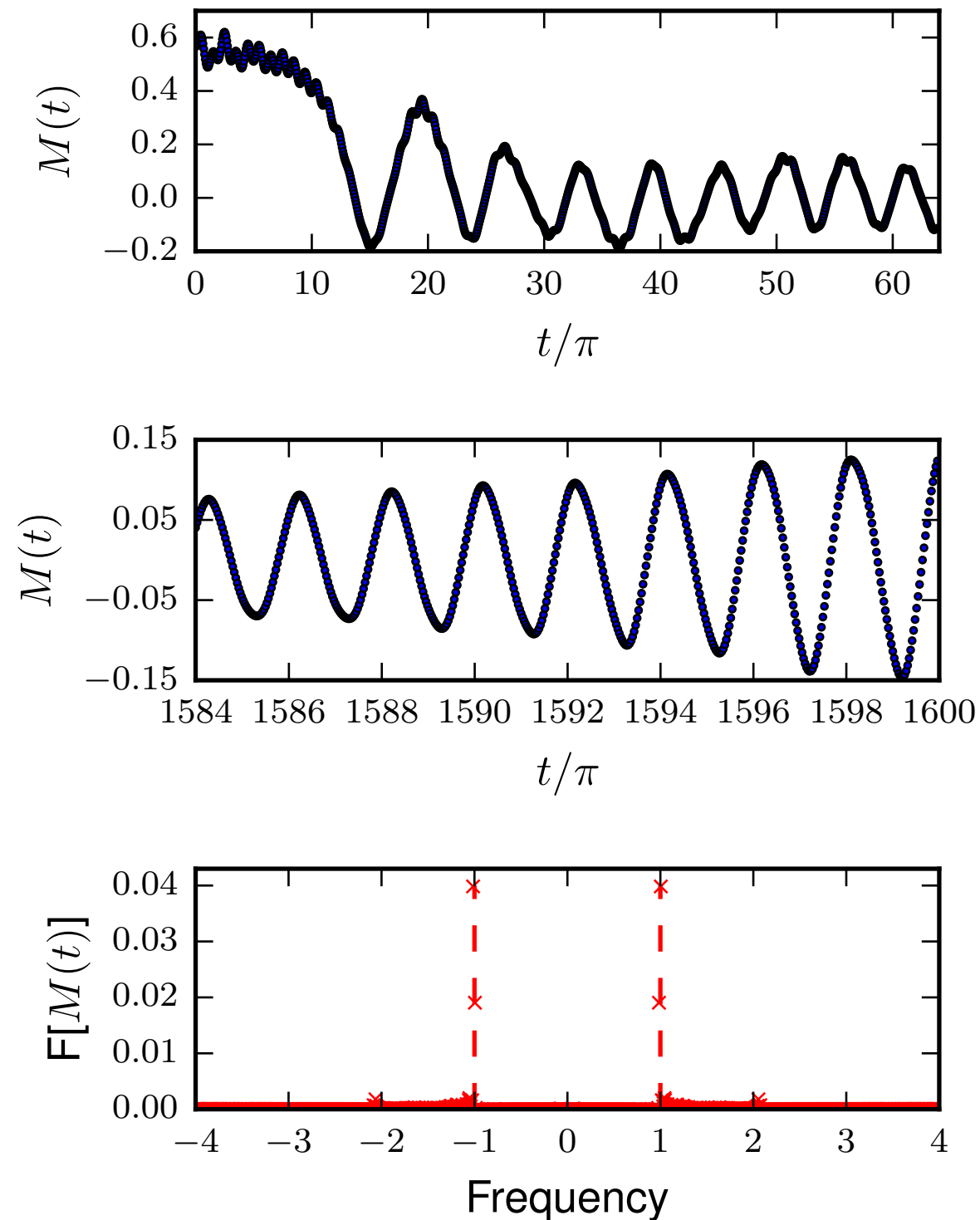
Feedback term prevents parametric resonance

Steady state: finite correlations with structure

“Integrable”: unknown generalized Gibbs ensemble

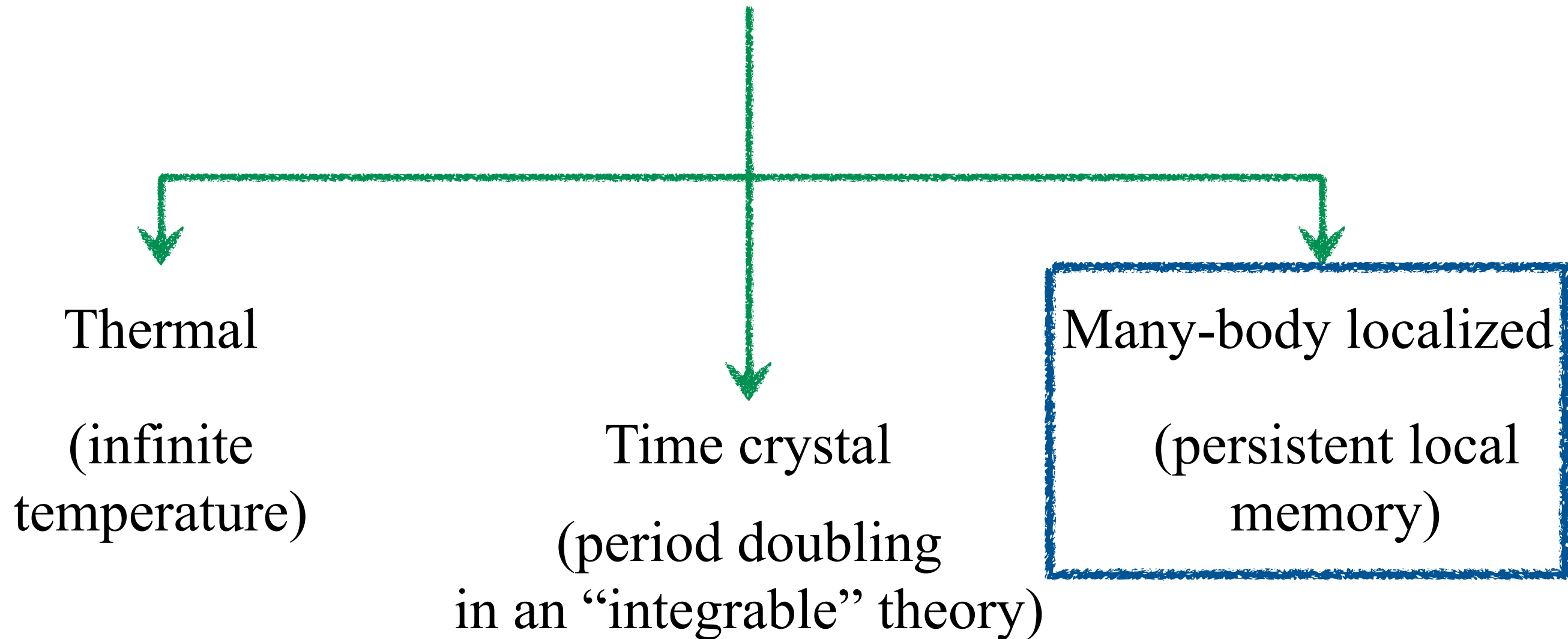
Period doubling in the driven ferromagnet

Drive period = π

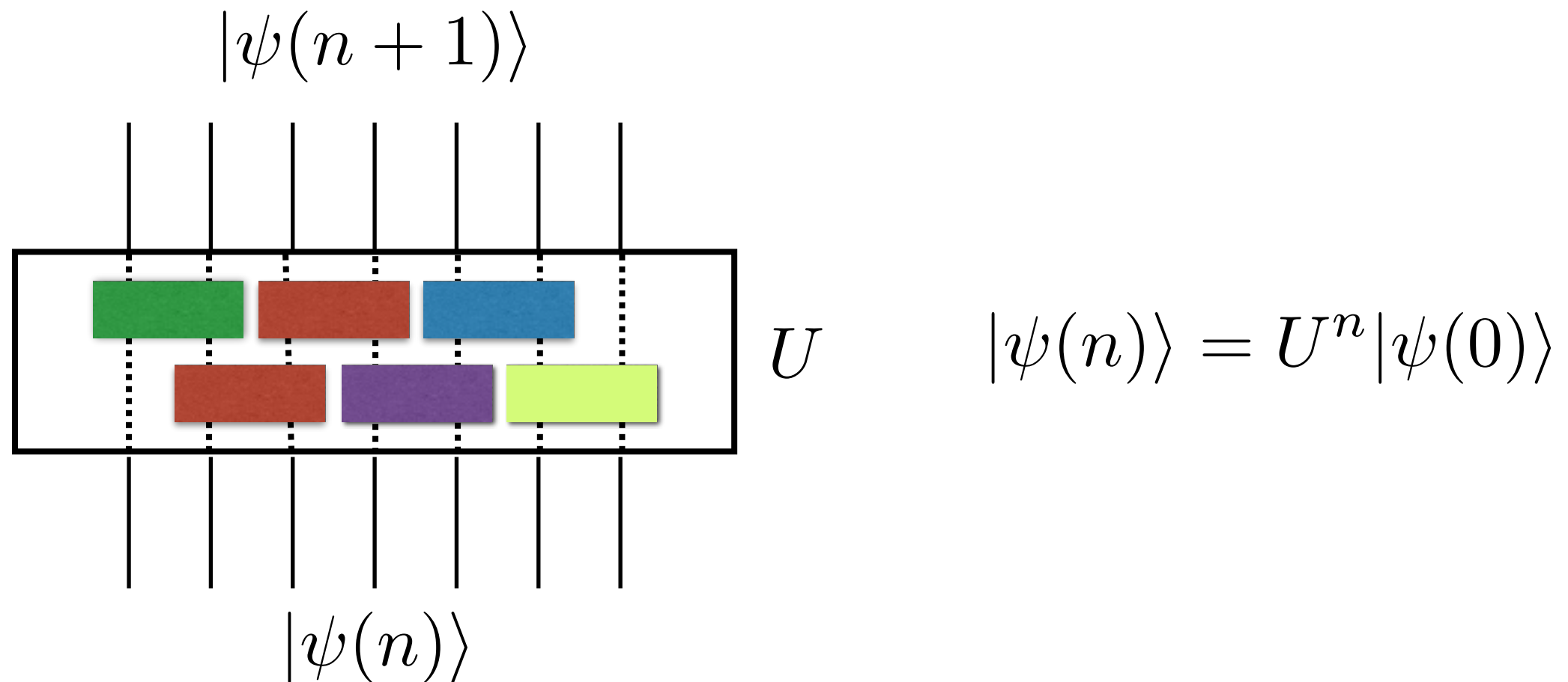


Outline

Steady states of Floquet systems



Periodic circuits

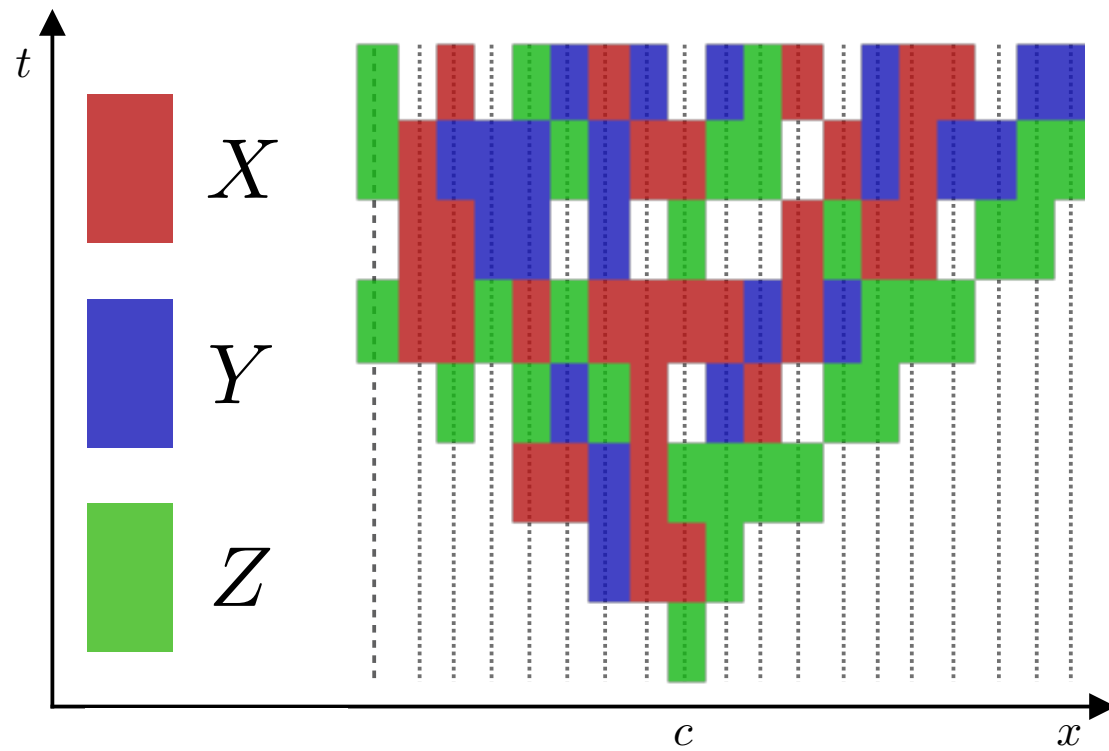


Floquet evolution without $H(t)$!

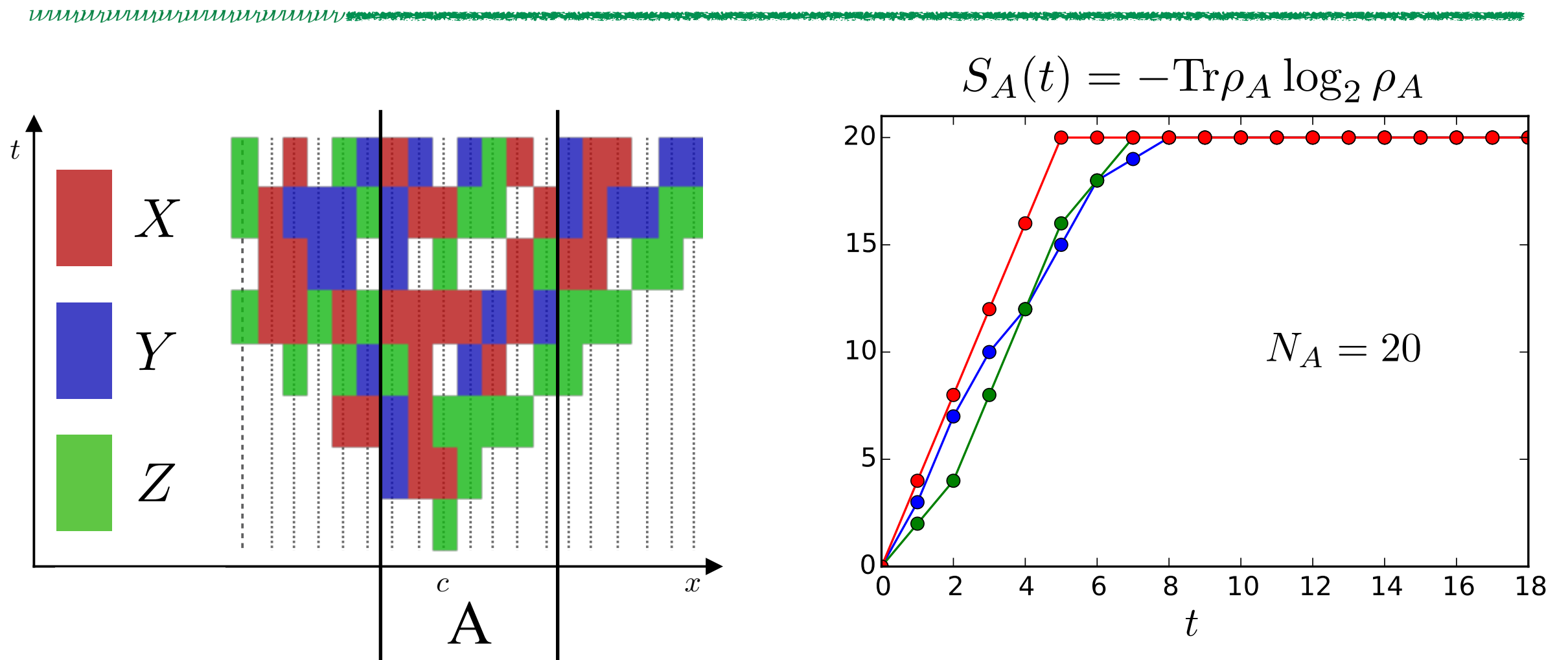
Clifford circuits

- Clifford gates: Hadamard, Phase, CNOT
- Efficiently simulable (poly(N) time for N qubits)
 - $U^\dagger (X_1 \otimes Z_2 \otimes \dots 1_N) U = Y_1 \otimes X_2 \otimes \dots Z_N$
- Can entangle
- Infinite temperature locally?

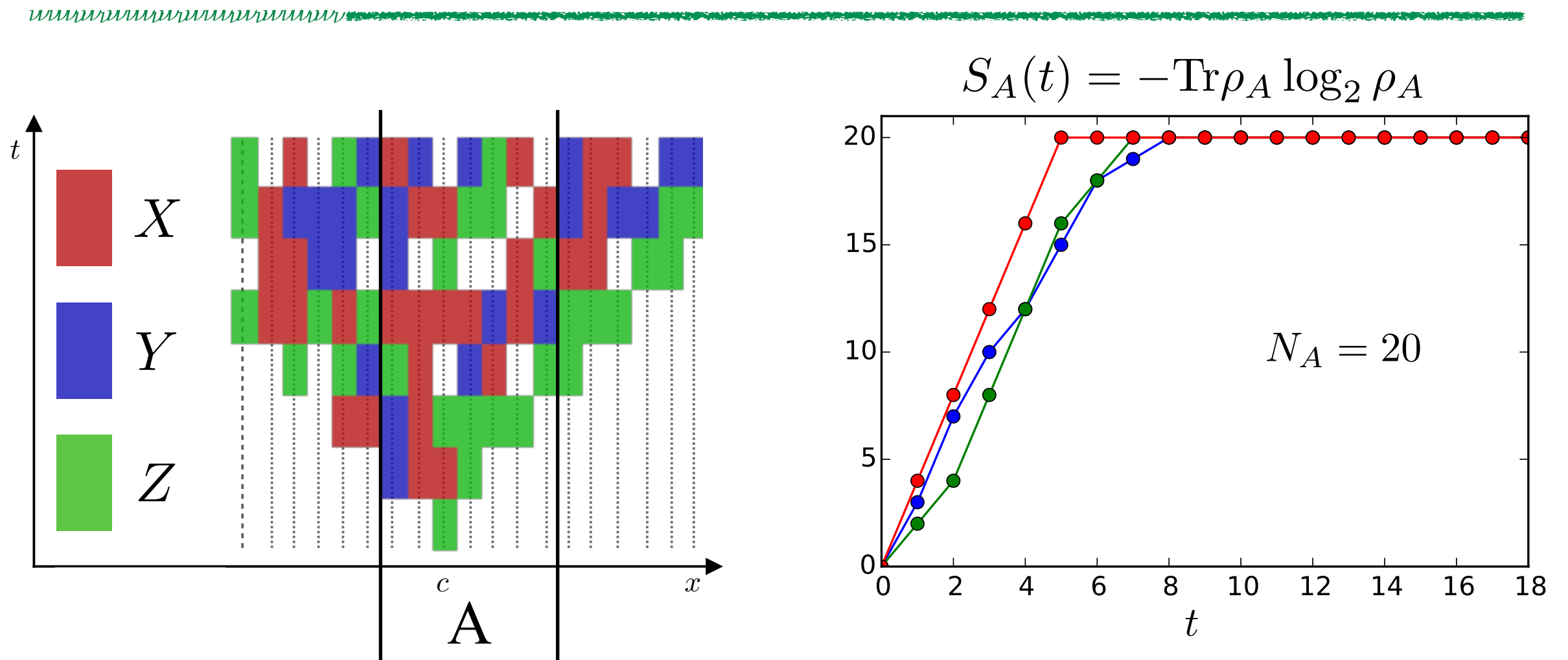
Thermalization



Thermalization

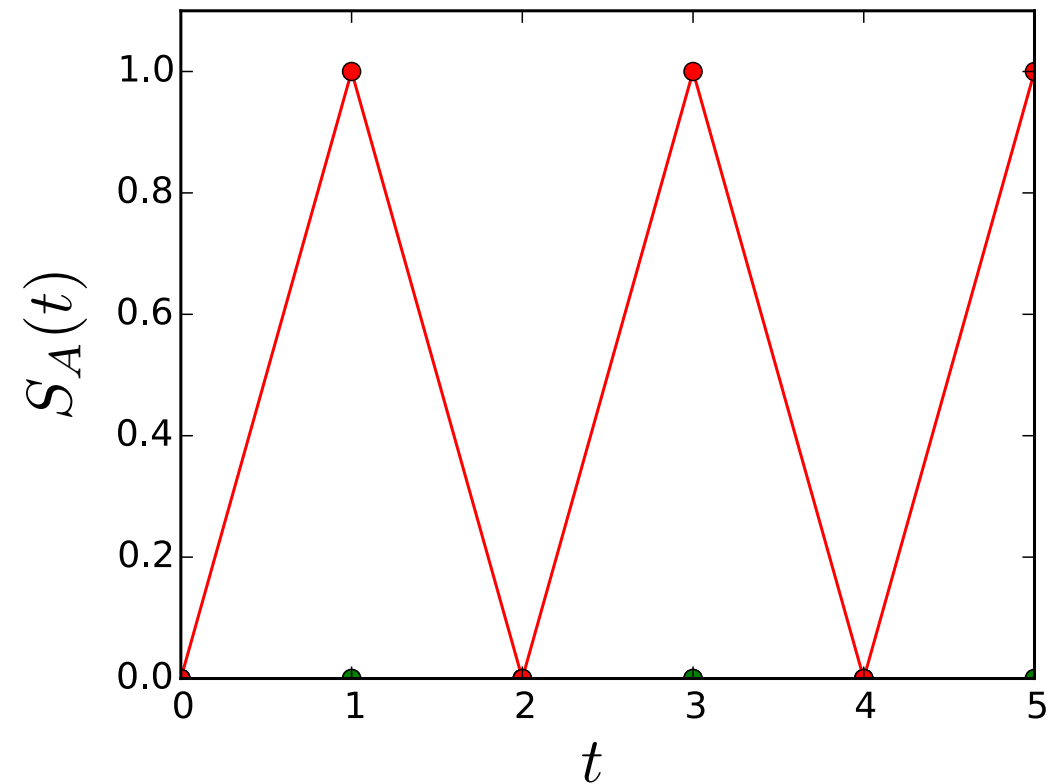
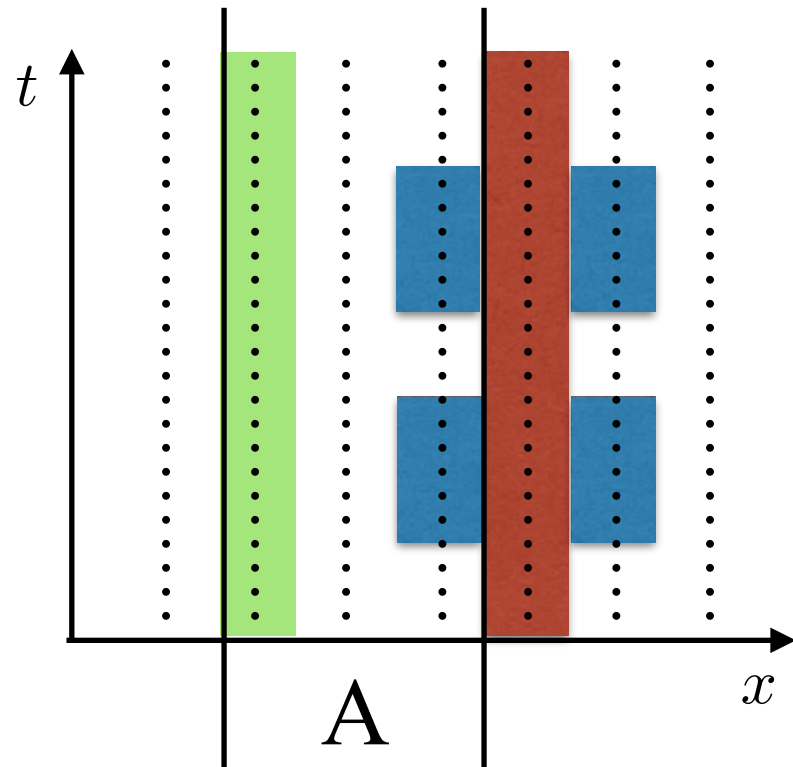


Thermalization



- Operator support grows in time
 - $\rho_A = 1$ for $t > vN_A$
- Simulable system that thermalizes!

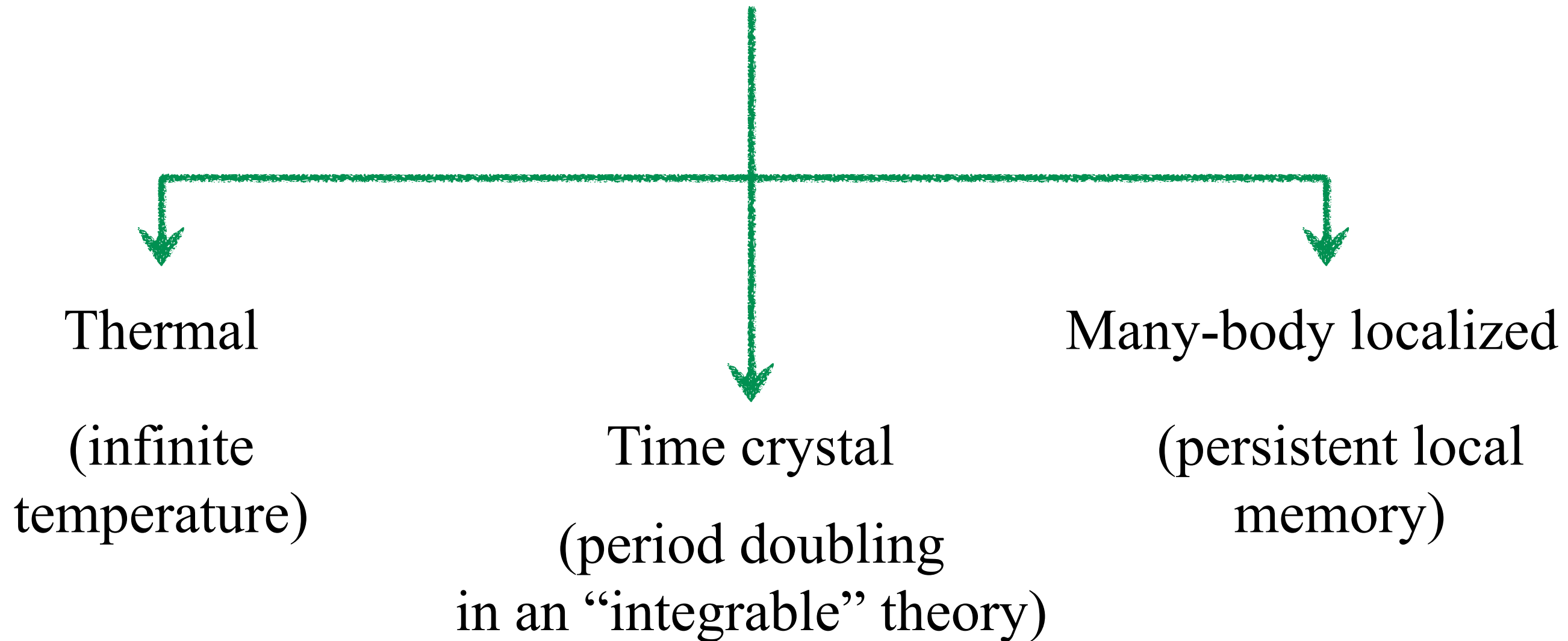
Localization



- Strictly local integrals of motion: Z_i
- Block spread of information
- Transition to thermalization: percolation of operator support

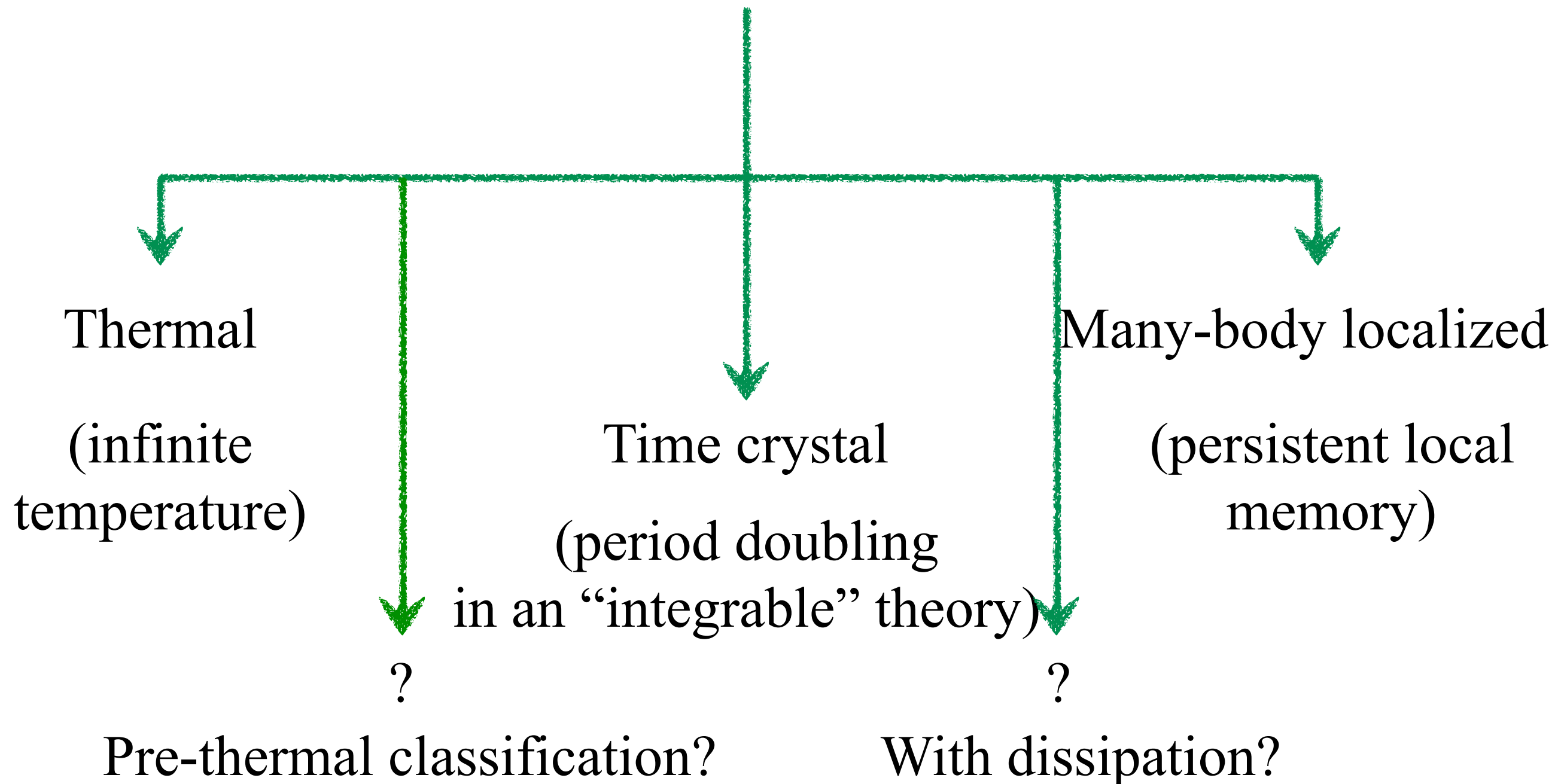
Outline

Steady states of Floquet systems



Outline

Steady states of Floquet systems



Thank you

Thank you to my collaborators

Dima Abanin, Chris Laumann,
Zlatko Papić, Pedro Ponte & Shivaji Sondhi

&

Thank you for your attention!

